



THE SOLUTION OF THE SECOND FUNDAMENTAL PROBLEM OF THE THEORY OF ELASTICITY FOR A PLATE WITH A DOUBLY SYMMETRIC TWO-CUSP CUT†

N. A. IVAN'SHIN and Ye. A. SHIROKOVA

Kazan'

(Received 15 May 1995)

As an extension of a previous paper [1], devoted to solving the first fundamental and mixed problems for a plate with a doubly symmetric two-cusp cut, the second fundamental problem for a plate with the same cut is solved by the same method. © 1997 Elsevier Science Ltd. All rights reserved.

1. THE SOLUTION PROCEDURE

The second boundary-value problem for an unbounded region D consists of finding two functions that are analytic in $D \setminus \{\infty\}$ [2]

$$f(z) = \Gamma z - \frac{X + iY}{2\pi(1 + \kappa)} \ln z + \frac{a}{z} + \dots, \quad g(z) = \Gamma' z + \frac{\kappa(X - iY)}{2\pi(1 + \kappa)} \ln z + \frac{a'}{z} + \dots \quad (1.1)$$

where Γ, Γ' and $X + iY$ are the known constants. The boundary condition has the form

$$\{\kappa f(z) - z \overline{f'(z)} - \overline{g(z)}\}|_{z=z(t)} = 2\mu(u(t) + iv(t)) \quad (1.2)$$

where $z = z(t), t \in [0, l]$ is the equation of the boundary curve ∂D . We will assume that $u'(t)$ and $v'(t)$ are Hölder's functions.

We will change to the function $z(\zeta)$, which conformally maps the region $E^- = \{\zeta = \xi + i\eta, |\zeta| > 1\}$ into D with corresponding $z(\infty) = \infty$ and differentiate both sides of equality (1.2) with respect to t . We obtain the boundary condition

$$\{\kappa \overline{\Phi(\zeta) z'(\zeta)} - \overline{z'(\zeta)} \Phi(\zeta) + z(\zeta) \overline{\Phi'(\zeta) \zeta^2} + \Psi(\zeta) z'(\zeta) \zeta^2\}|_{\zeta=e^{i\theta}} = Q(\theta) \quad (1.3)$$

$$Q(\theta) = 2i\mu[u'(t) - iv'(t)]|_{t=t(\theta)} e^{i\theta} |z'(e^{i\theta})|, \quad \theta \in [-\pi, \pi] \quad (1.4)$$

where the functions $\Phi(\zeta) = f'(z(\zeta))$ and $\Psi(\zeta) = g'(z(\zeta))$ are analytic in E^- .

If $q(\zeta) = z^-(\zeta^{-1})$ and $r(\zeta) = z^-(\zeta^{-1})$ are meromorphic in E^- , we can restore the meromorphic functions

$$K_j(\zeta) = (-1)^j \kappa \Phi(\zeta) z'(\zeta) - r(\zeta) \Phi(\zeta) + q(\zeta) \Phi'(\zeta) \zeta^2 + \Psi(\zeta) z'(\zeta) \zeta^2, \quad j = 1, 2 \quad (1.5)$$

using the boundary values $\text{Re } K_1(\zeta)$ ($\text{Re } Q(\theta)$) and $\text{Im } K_2(\zeta)$ ($\text{Im } Q(\theta)$), (1.3) and (1.4). Further, $\Phi(\zeta)$ and $\Psi(\zeta)$ are restored by $K_j(\zeta)$ (1.5).

2. SOLUTION OF THE PROBLEM

The function

$$z(\zeta) = \frac{i(b^2 \zeta^2 + 1)}{\zeta(b^2 - 1)} + \frac{\zeta(b^2 - 1)}{i(b^2 \zeta^2 + 1)}, \quad b > 1 \quad (2.1)$$

maps E^- into D . Using Eq. (1.5), we obtain

$$2\kappa \Phi(\zeta) z'(\zeta) = \frac{1}{2\pi} \int_{-\pi}^{\pi} \overline{Q(\theta)} \frac{\zeta + e^{i\theta}}{\zeta - e^{i\theta}} d\theta + A - \frac{2}{\zeta^2} \left[\frac{i}{b^2 - 1} \overline{\Gamma} - \frac{ib^2}{b^2 - 1} \overline{\Gamma'} \right] -$$

†Prikl. Mat. Mekh. Vol. 61, No. 2, pp. 350-351, 1997.

$$-\frac{\kappa(X+iY)}{\pi(1+\kappa)\zeta} - 2\bar{B}\frac{\zeta-ib}{1+ib\zeta} - 2\bar{C}\frac{\zeta+ib}{1-ib\zeta} - 2\bar{D}\left(\frac{\zeta-ib}{1+ib\zeta}\right)^2 - 2\bar{E}\left(\frac{\zeta+ib}{1-ib\zeta}\right)^2 \quad (2.2)$$

Requiring that $\Phi(\infty) = \Gamma$, we have

$$A + \frac{2i}{b}\bar{B} - \frac{2i}{b}\bar{C} + \frac{2}{b^2}\bar{D} + \frac{2}{b^2}\bar{E} = \frac{2\kappa ib^2}{b^2-1}\Gamma - \frac{1}{2\pi} \int_{-\pi}^{\pi} \overline{Q(\theta)} d\theta \quad (2.3)$$

By (1.5)

$$\Psi(\zeta) = [\zeta^2 z'(\zeta)]^{-1} [K_1(\zeta) - \kappa\Phi(\zeta)z'(\zeta) + r(\zeta)\Phi(\zeta) - q(\zeta)\Phi'(\zeta)\zeta^2] \quad (2.4)$$

where the coefficients of $(\zeta \pm ib)^{-2}$ and $(\zeta + ib)^{-1}$ must vanish, which yields four equations

$$\begin{aligned} D(b^2-1) - i\Phi(ib)b^2/2 &= 0, & E(b^2-1) - i\Phi(-ib)b^2/2 &= 0 \\ B + 2ibD + b\Phi(ib) &= 0, & C - 2ibE - b\Phi(-ib) &= 0 \end{aligned} \quad (2.5)$$

Using formula (2.2), we obtain from (2.3) and (2.5) a system of five equations in five unknown constants

$$\begin{aligned} A + \frac{2i}{b}\bar{B} - \frac{2i}{b}\bar{C} + \frac{2}{b^2}\bar{D} + \frac{2}{b^2}\bar{E} &= \frac{2\kappa ib^2}{b^2-1}\Gamma - \frac{1}{2\pi} \int_{-\pi}^{\pi} \overline{Q(\theta)} d\theta \\ D - b\beta[A - 2i\delta\bar{C} + 2\delta^2\bar{E}] &= b\beta R_+(b) \\ E - b\beta[A + 2i\delta\bar{B} + 2\delta^2\bar{D}] &= b\beta R_-(b) \\ B + 2ibD + 2i\beta(1-b^2)[A - 2i\delta\bar{C} + 2\delta^2\bar{E}] &= -2i\beta(1-b^2)R_+(b) \\ C - 2ibE - 2i\beta(1-b^2)[A + 2i\delta\bar{B} + 2\delta^2\bar{D}] &= 2i\beta(1-b^2)R_-(b) \end{aligned} \quad (2.6)$$

Here

$$\begin{aligned} \beta &= \frac{b^3(1+b^2)^2}{4\kappa(b^4+1)(b^4+b^2+1)}, & \delta &= \frac{2b}{b^2+1} \\ R_{\pm}(b) &= \frac{1}{2\pi} \int_{-\pi}^{\pi} \overline{Q(\theta)} \frac{\pm ib + e^{i\theta}}{\pm ib - e^{i\theta}} d\theta - \frac{2i}{b^2(b^2-1)}(\bar{\Gamma} - b^2\bar{\Gamma}') \mp \frac{\kappa(X+iY)}{\pi ib(1+\kappa)} \end{aligned}$$

After determining the constants, the functions $\Phi(\zeta)$ and $\Psi(\zeta)$ can be obtained from (2.2) and (2.4) respectively. Now we reconstruct $f(z(\zeta))$ and $g(z(\zeta))$

$$f(z(\zeta)) = \int_1^{\zeta} \Phi(\zeta)z'(\zeta)d\zeta + f(z(1)), \quad g(z(\zeta)) = \int_1^{\zeta} \Psi(\zeta)z'(\zeta)d\zeta + g(z(1))$$

The function $z(\zeta)$ is the same as in (2.1), $f(z(1))$ and $g(z(1))$ are the quantities that satisfy the equation

$$\kappa f(z(1)) - \overline{g(z(1))} = 2\mu[u(t_0) + iv(t_0)] + z(1)\overline{\Phi(1)}$$

and t_0 is the value of the initial parameter corresponding to the point $z = z(1)$.

REFERENCES

1. IVAN'SHIN, N. A. and SHIROKOVA, Ye. A., The solution of the problem of the theory of elasticity for a plane with a doubly symmetric two-cusp cut. *Prikl. Mat. Mekh.*, 1995, 59, 3, 497-501.
2. MUSKHELISHVILI, N. I., *Some Fundamental Problems of the Mathematical Theory of Elasticity*. Nauka, Moscow, 1976.

Translated by N.R.